

6. 自由電子ガス・平面波基底

6.1 一電子軌道と k 点種分

$$\left\{ -\frac{\nabla^2}{2} + \cancel{U(r)} \right\} \varphi_{k_j}(r) = \varepsilon_{k_j} \varphi_{k_j}(r)$$

$\cancel{U(r)}$ は定数

$$-\frac{\nabla^2}{2} \varphi_k(r) = \varepsilon_k \varphi_k(r) \Rightarrow \varphi_k = \frac{e^{ikr}}{\sqrt{V}}, \quad \varepsilon_k = \frac{k^2}{2}$$

$$u_k = \frac{1}{\sqrt{V_{uc}}} : \text{定数}$$

$$\hookrightarrow N_{uc} V_{uc}$$

$$\sum_k^{N_{uc}} \dots \Rightarrow \int_{BZ} d^3k \frac{N_{uc}}{V_{BZ}} \dots = \frac{V}{(2\pi)^3} \int_{BZ} d^3k \dots$$

$$\hookrightarrow \|(b_1, b_2, b_3)\| = (2\pi)^3 \|(a_1, a_2, a_3)^{t^{-1}}\| = \frac{(2\pi)^3}{V_{uc}}$$

自由電子系では、 $V_{uc} \rightarrow 0, V_{BZ} \rightarrow \infty$ となる

$$\sum_k \dots \rightarrow \frac{V}{(2\pi)^3} \int_0^\infty dk k^2 \int_0^\pi d\theta \sin\theta \int_0^{2\pi} d\phi \dots$$

6.2 物理量と DOS

・エネルギー (単位体積あたり)

$$\frac{E}{V} = \frac{1}{(2\pi)^3} \int_0^\infty dk k^2 \int_0^\pi d\theta \sin\theta \int_0^{2\pi} d\phi \frac{k^2}{2} \theta\left(\varepsilon_F - \frac{k^2}{2}\right)$$

$$\frac{D(\varepsilon)}{V} = \frac{1}{(2\pi)^3} \int_0^\infty dk k^2 \int_0^\pi d\theta \sin\theta \int_0^{2\pi} d\phi \delta\left(\varepsilon - \frac{k^2}{2}\right)$$

レポート課題

6.3 介極関数

$$\Pi_0^{(q)}(r, r') \equiv \sum_{jj'} \int_{BZ} \frac{d^3k}{V_{BZ}} \frac{-f_{kj}(1-f_{k+qj'})}{\epsilon_{k+qj'} - \epsilon_{kj}} \rho_{kj, k+qj'}(r) \rho_{kj, k+qj'}(r')$$

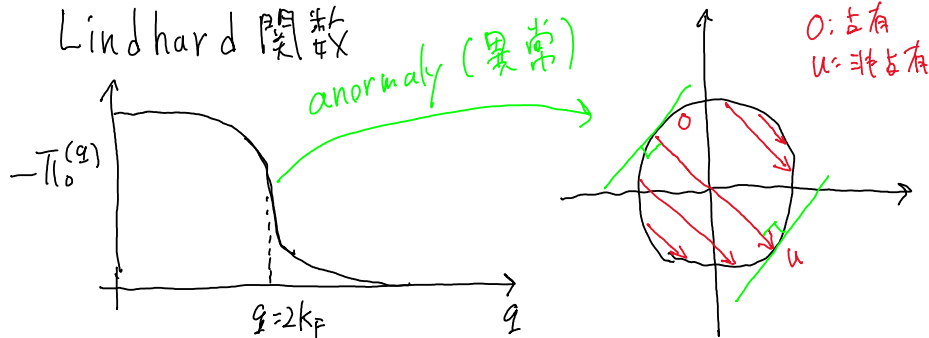
↓ r, r' 非依存

$$u_k = \frac{1}{\sqrt{V_{uc}}} \quad \tau_j \text{ の } 2^{\text{nd}} \quad \frac{1}{V_{uc}}$$

$$\Pi_0^{(q)} = \frac{1}{V_{uc}} \int_{BZ} \frac{d^3k}{(2\pi)^3} \frac{-f_k(1-f_{k+q})}{\epsilon_{k+q} - \epsilon_k}$$

⇓

Lindhard 関数



6.4 平面波基底

一般のポテンシャル $v(r)$

$$\left(-\frac{(\nabla + ik)^2}{2} + v(r) \right) u_{kj}(r) = \epsilon u_{kj}(r) \quad \text{を解く (数値的)}$$

$$u_{kj}(r) = \sum_{\alpha} \tilde{u}_{kj}(\alpha) \chi_{\alpha}(r) \quad \text{基底展開}$$

↳ $\chi_{\alpha}(r) = \delta(r - r_{\alpha}) \rightarrow$ 実空間 χ_{α} シュ
有限要素法

平面波展開

$$u_{kj}(r) = \sum_{G} e^{iG \cdot r} \tilde{u}_{kj}(G)$$

“ 逆格子ベクトル $G = n_1 b_1 + n_2 b_2 + n_3 b_3$

$$\left(\frac{(G+k)^2}{2} + \hat{V} \right) \tilde{u}_{kj}(G) = \epsilon \tilde{u}_{kj}(G)$$

$\downarrow$
 $V(G, G')$ のような $\tilde{u}_{kj}(G')$ に $\tilde{u}_{kj}(G)$ になる...

$$\hat{V} \tilde{u}_{kj}(G) \equiv \sum_{G'} V(G, G') \tilde{u}_{kj}(G') = \int_{uc} d^3r \frac{e^{-iGr}}{\sqrt{V_{uc}}} v(r) u_{kj}(r)$$

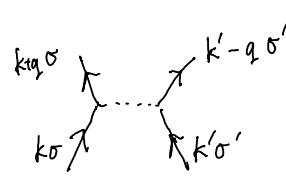
$\swarrow$
 $O(N_{pw}^2)$

$\downarrow$
 $T = T_0 + \frac{1}{N_{pw}} \rightarrow$ 高速フーリエ変換
 $O(N_{pw} \ln N_{pw})$

6.5 自由電子の生成・消滅演算子

$$\hat{\psi}_\sigma(r) = \sum_k \hat{C}_{k\sigma} \frac{e^{ikr}}{\sqrt{V}} = \frac{\sqrt{V}}{(2\pi)^3} \int d^3k C_{k\sigma} e^{ikr}$$

$$\begin{aligned} H^{(1)} &= \sum_\sigma \sum_{kk'} C_{k\sigma}^\dagger C_{k'\sigma} \int_V d^3r \frac{1}{V} e^{-ikr} \left\{ -\frac{\nabla^2}{2} + v(r) \right\} e^{ik'r} \\ &= \sum_\sigma \sum_{kk'} C_{k\sigma}^\dagger C_{k'\sigma} \left\{ \underbrace{\frac{k'^2}{2} \int_V \frac{d^3r}{V} e^{i(k'-k)r}}_{\delta_{kk'}} + \underbrace{\frac{1}{V} \int_V d^3r v(r) e^{i(k'-k)r}}_{v_{kk'}} \right\} \\ &= \sum_\sigma \left(\sum_k \frac{k^2}{2} C_{k\sigma}^\dagger C_{k\sigma} + \sum_{kk'} v_{kk'} C_{k\sigma}^\dagger C_{k'\sigma} \right) \end{aligned}$$

$$\begin{aligned}
 H^{(2)} &= \frac{1}{2} \sum_{\sigma\sigma'} \sum_{k_1 k_2 k_3 k_4} C_{k_1\sigma}^\dagger C_{k_2\sigma'}^\dagger C_{k_3\sigma'} C_{k_4\sigma} \frac{1}{V^2} \int_V d^3r \int_V d^3r' \frac{e^{i(k_4-k_1)r} e^{i(k_3-k_2)r'}}{|r-r'|} \\
 &\quad \xrightarrow{\substack{r-r' \rightarrow r \\ r \rightarrow r+r'}} \frac{1}{V^2} \int d^3r \frac{e^{i(k_4-k_1)r}}{|r|} \int d^3r' e^{i(k_3+k_4-k_1-k_2)r'} \xrightarrow{\substack{k_3 \rightarrow k' \\ k_4 \rightarrow k \\ k_1 \rightarrow k+q \\ k_2 \rightarrow k'-q}} \\
 &\quad \rightarrow 2\pi \int_0^\infty dr r^2 \int_{-1}^1 dt \frac{e^{ikrt}}{r} = \frac{4\pi}{k^2} \\
 &= \frac{1}{2} \sum_{\sigma\sigma'} \sum_{kk'q} \frac{4\pi}{q^2} C_{k+q\sigma}^\dagger C_{k'-q\sigma'}^\dagger C_{k'\sigma'} C_{k\sigma}
 \end{aligned}$$


• 反交換関係

$$\begin{aligned}
 C_{k\sigma} &= \frac{1}{\sqrt{V}} \int d^3r e^{-ikr} \psi_\sigma(r) \\
 \{C_{k\sigma}, C_{k'\sigma'}^\dagger\} &= \frac{1}{V} \int d^3r \int d^3r' e^{-ikr} e^{ik'r'} \underbrace{\{\psi_\sigma(r), \psi_{\sigma'}^\dagger(r')\}}_{\delta_{\sigma\sigma'} \delta(r-r')} \\
 &= \frac{1}{V} \int d^3r e^{i(k'-k)r} \delta_{\sigma\sigma'} = \delta_{kk'} \delta_{\sigma\sigma'}
 \end{aligned}$$